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Simulations of non-local spin interaction in atomic magnetometers Title:

using LANL's D-Wave 2X

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Simulations of non-local spin interaction in atomic magnetometers using LANL's D-Wave 2X

Igor Savukov and Alexander Malyzhenkov

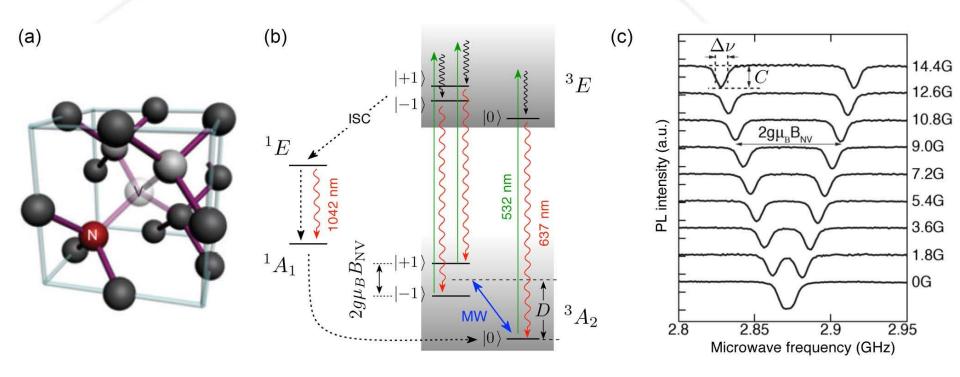
4/27/2017



Nitrogen-vacancy (NV) defect in diamond



$$H = hDS_z^2 + hE(S_x^2 - S_y^2) + g\mu_B BS + H_d$$



[1] Rep. Prog. Phys. **77** (2014) 056503 (26pp)

National Nuclear Security Administration
Slide 2





D ~ 2.87 GHz or 100 mT

E ~ 100 kHz and can be neglected

 $H_d \sim 100 \text{ mT for } 1 \text{ nm}, \text{ or } 0.1 \text{ mT for } 10 \text{ nm}.$

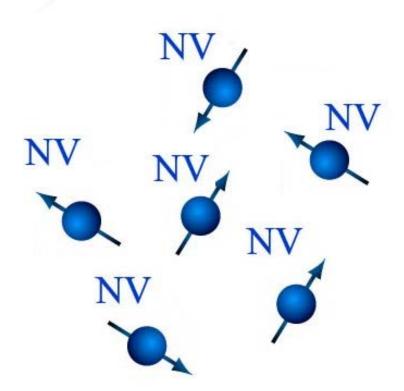
 hDS_z^2 is not sensitive to the direction and does not play any role in our analysis when we consider only z component; however, forces the spin to be parallel to z.



Ensemble of NV defects:



Signal to noise ratio for independent NV-centers:



$$SNR \sim \frac{1}{\sqrt{N}}$$

But dipole interaction between different NVcenters will also impact sensitivity and orientation of spins

$$H_i = -\sum \boldsymbol{m_{NV}}_i \boldsymbol{B_{NV}}_j$$

$$\mathbf{B}(\mathbf{r}) =
abla imes \mathbf{A} = rac{\mu_0}{4\pi} \left(rac{3\mathbf{r}(\mathbf{m} \cdot \mathbf{r})}{r^5} - rac{\mathbf{m}}{r^3}
ight)$$







Spin-spin interaction with random positioning is extremely hard problem for classical computer, but it seems to be an ideal fit for Ising Quantum Computer such as D-Wave

Approximation: spin and magnetic field only along/opposite with z-axis; however, it can be justified for stationary states

$$H = \sum_{i} k_{1} s_{i} B(r_{i}) +$$

$$\sum_{i} \sum_{j < i} 2k_{2} s_{i} s_{j} / r_{ij}^{3} (1 - \frac{3z_{ij}^{2}}{r_{ij}^{2}})$$

$$r_{ij} = \sqrt{(i_x - j_x)^2 + (i_y - j_y)^2 + (i_z - j_z)^2}; z_{ij} = (i_z - j_z)$$



Ground and excited states



To understand the dynamics of NV-diamond spins, the ground state is not sufficient and excited states also need to be found.

D-Wave finds automatically the ground state or low-excited states when the temperature is finite and the thermal energy is comparable to the splitting.

A more reliable method to find excited states is to introduce penalty energy for a state of being non-orthogonal to the ground state

$$H = k_3 \sum_i s_i^{gr} s_i, k_3 > 0$$

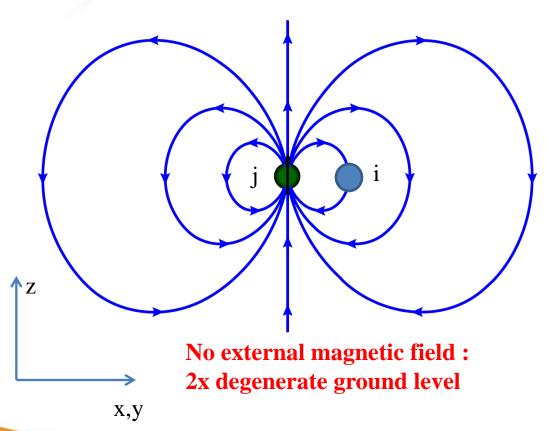
This term can be combined with the Zeeman single-sum term and the new "ground" state, orthogonal to the real one is found by D-Wave. This process can be continued by adding the term enforcing orthogonality with the first excited states, and so on.

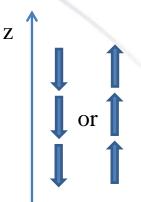


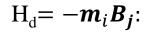
Dipole field map to predict solutions



We can predict solutions in some trivial spin configurations



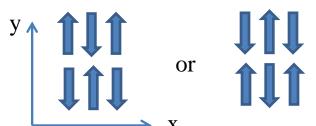




Spin is aligned with magnetic field of its closest neighbor to minimize Hamiltonian





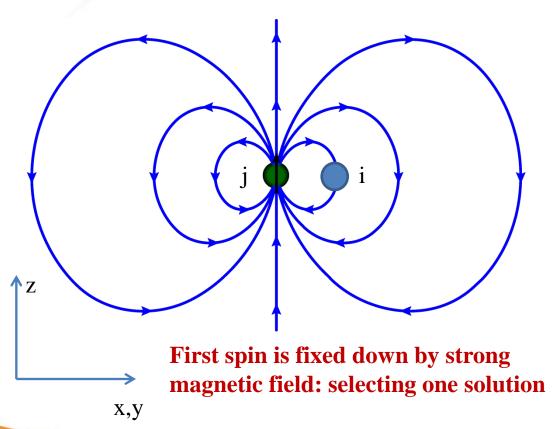


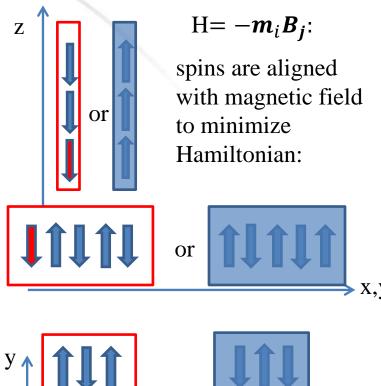


Dipole field map to predict solutions



We can predict solutions in some trivial spin configurations





or





Is D-Wave 2X capable to find energy levels (at least ground state) of many (how many?) NV-diamond spins system in our approximation?





Problems we found

1. D-Wave 2X bias: spin up and spin down are not the same

2. Embedding:

a) increasing the number of spins in the system requires a lot of couplings, which squeezes the dynamic range b) only 45 spins connected to each other are possible to simulate on D-Wave 2X unless we "cheat" and cut long distance interaction



2-spin system (x/y-line): no embedding needed



D-Wave solutions:

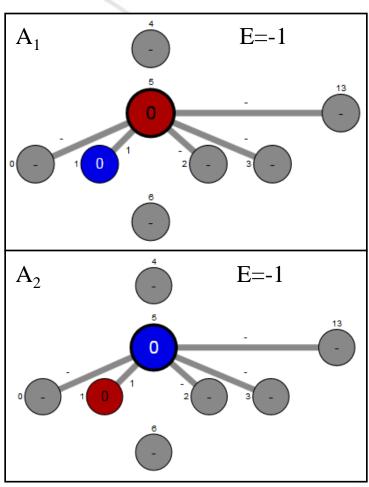
Expectations: 50% and 50%



Samples=1000

Run	1	2	3	4	5	Total
A_1	565	570	474	612	440	2661
A_2	435	430	526	388	560	2339

Classical computer *Quick sort: states to compute E $D_{min} = n \log(n)$ and sort*: n=2²=4 $D_{max} = n^2$

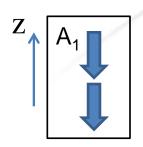


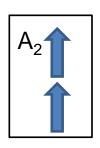


2-spin system (z-line): no embedding needed



Expectations: 50% and 50%





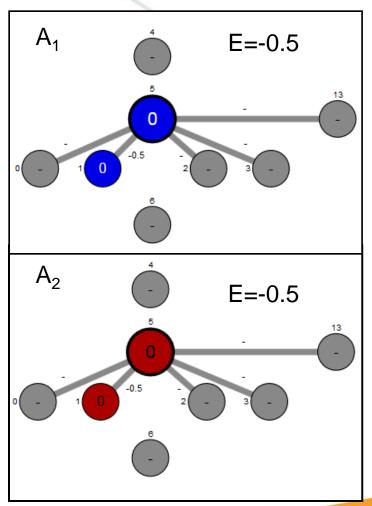
Samples=1000;



1

Run	1	2	3	4	5	Total
A_1	714	626	683	755	658	3436
A_2	286	374	317	245	342	1564

D-Wave solutions:



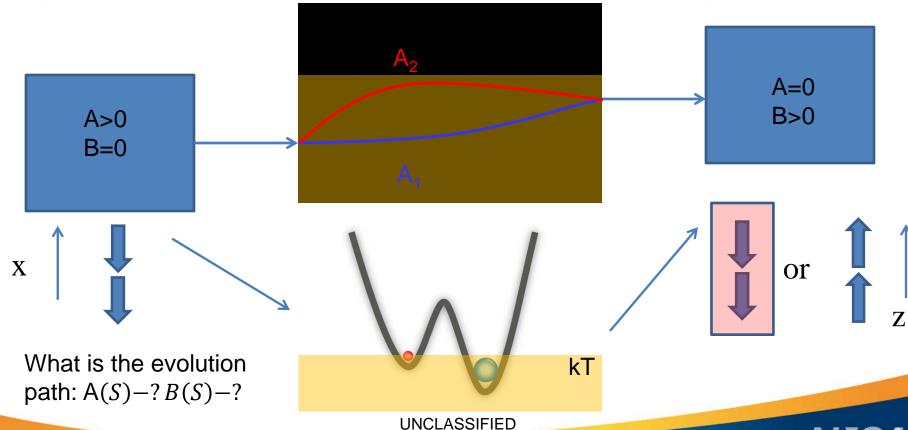


Bias: adiabatic DWAVE evolution



$$E_{ising}(s) = \frac{A(s)}{2} \left(\sum_{i} \sigma_x^{(i)} \right) + \frac{B(s)}{2} \left(\sum_{i} h_i \sigma_z^{(i)} + \sum_{i>j} J_{i,j} \sigma_z^{(i)} \sigma_z^{(j)} \right)$$

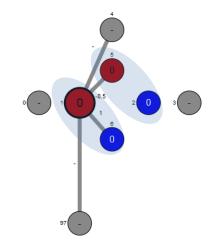
Bias dependence from temperature (annealing time) will be an interesting problem to study



How to get right solutions with biased D-Wave

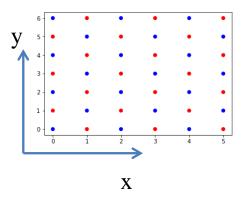


1. Spin anti-cloning *Problem: squeezing dynamic range* *



Anti-cloning based embedding would be ideal for systems which need embedding anyway *dynamic range would be the same as in regular embedding

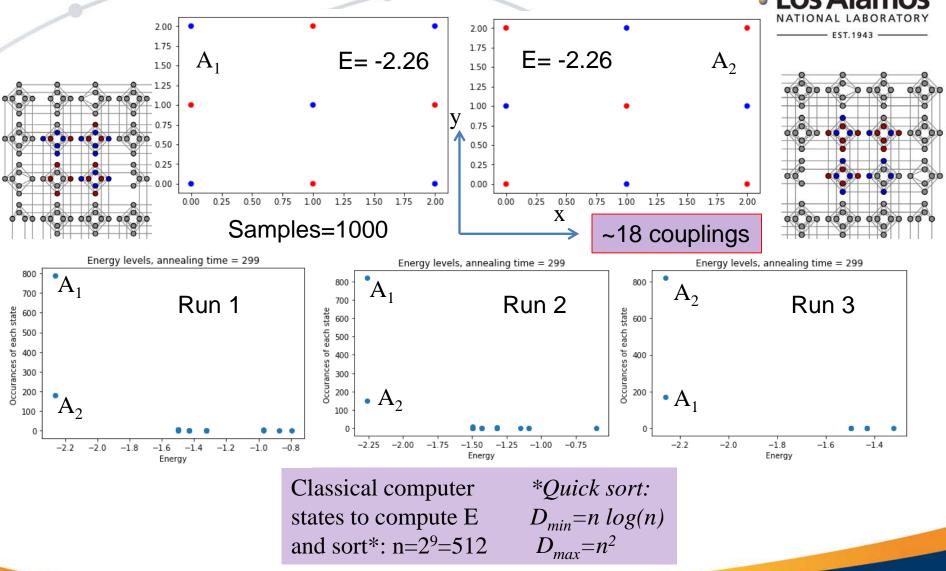
2. Choosing symmetrical relatively to spin direction systems to study



For such symmetrical systems all embedding chains should be ideally of the same length



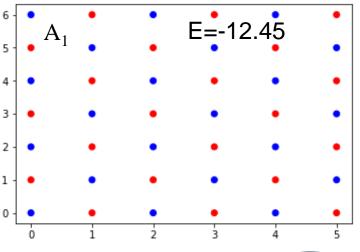
2D spin system: 9=3*3 spins

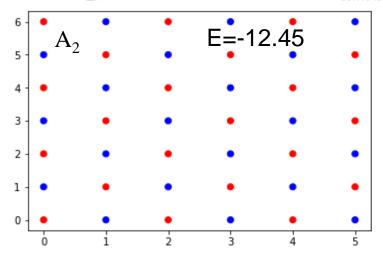


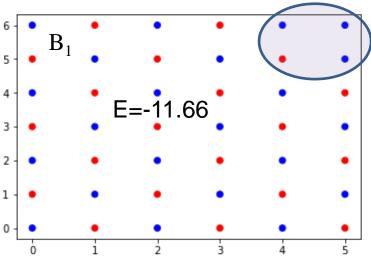


2D spin system: 42=6*7 spins









Spin flip

~800 couplings

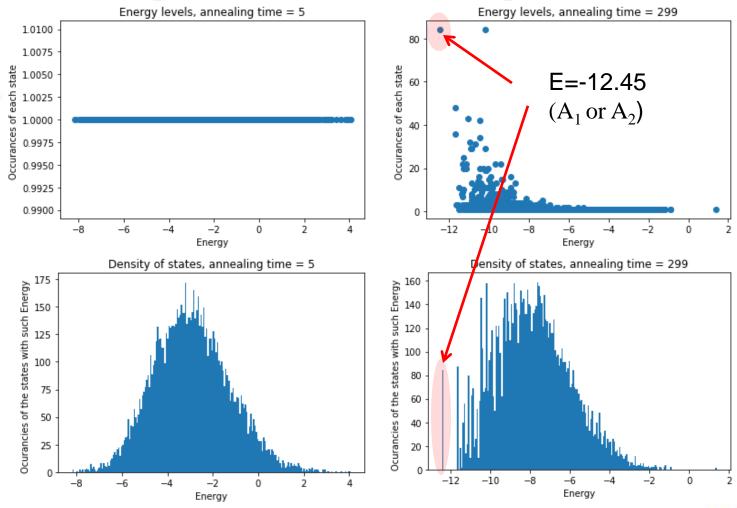
Classical computer states to compute E and sort*: $n=2^{42} \sim 4*10^{12}$ *Quick sort: $D_{min}=n log(n)$ $D_{max}=n^2$



Annealing time dependence

42 spins, samples=10000





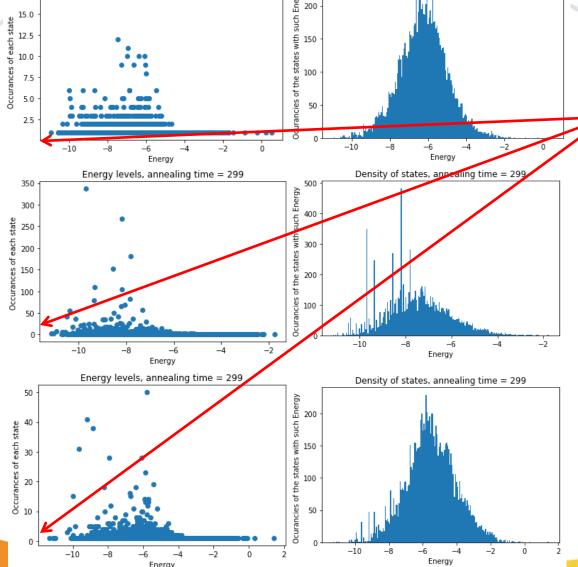


Same parameters – different results

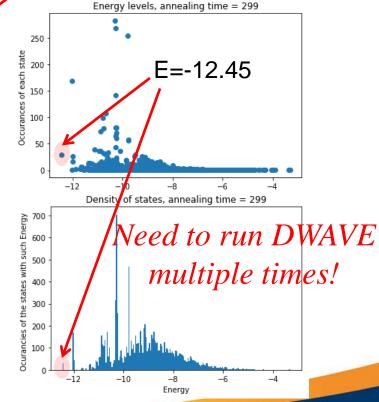
Density of states, annealing time = 299



Missing ground level solution 3 times from 4



Energy levels, annealing time = 299





Annealing time dependence: multiple runs

Only analyzing two first (lowest energy) levels from DWAVE each run:

[runs=400, samples=10000]

Solution 0 and 1 should have

the same energy ~-12.45, but we find only "0" in ground state

